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Phase mixing with multiple harmonics

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Introduction

- Our aim is to investigate laminar phase mixing of Alfvén waves as a standalone heating mechanism in coronal loops.
- We calculate an upper bound for the damping rate and compare with the heating requirement of the corona.
- We define the damping rate, γ , as

$$\gamma = \frac{\langle \text{Steady-state heating rate} \rangle}{\langle 2 \times \text{Steady-state kinetic wave energy} \rangle}$$

Model

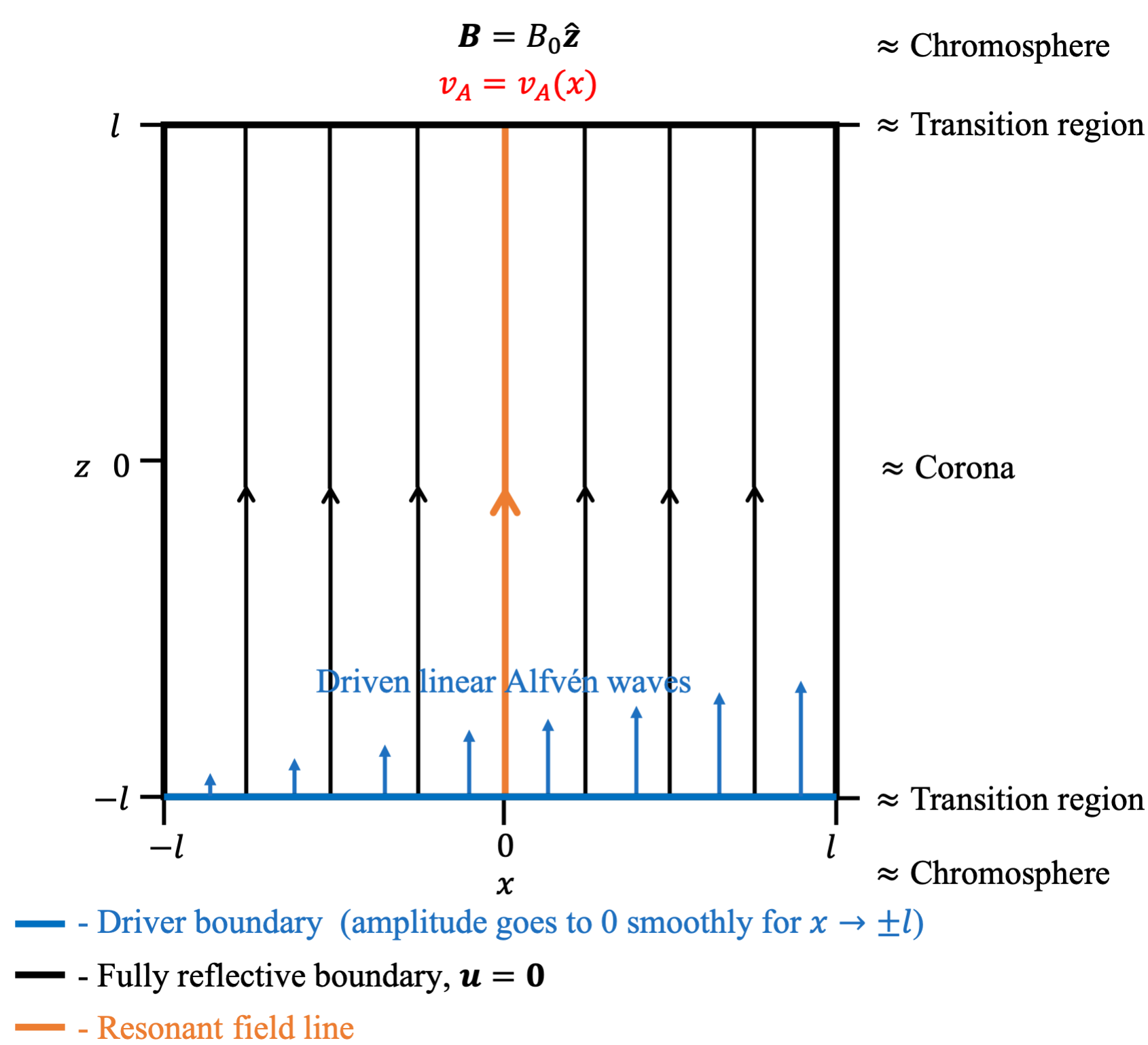


Fig. 1. Diagram of our model.

- **Driver** of the form:

$$\mathbf{u} = u_0 \sum_{n=1}^N A_n (n\omega_1)^{-\alpha/2} \sin(n\omega_1 t + \phi_n) \hat{\mathbf{y}},$$

ω_1 is the fundamental angular frequency of the resonant field line, ϕ_n is a random phase.

- Steady-state wave energy power spectrum with slope $-\alpha$ (Fig. 2).

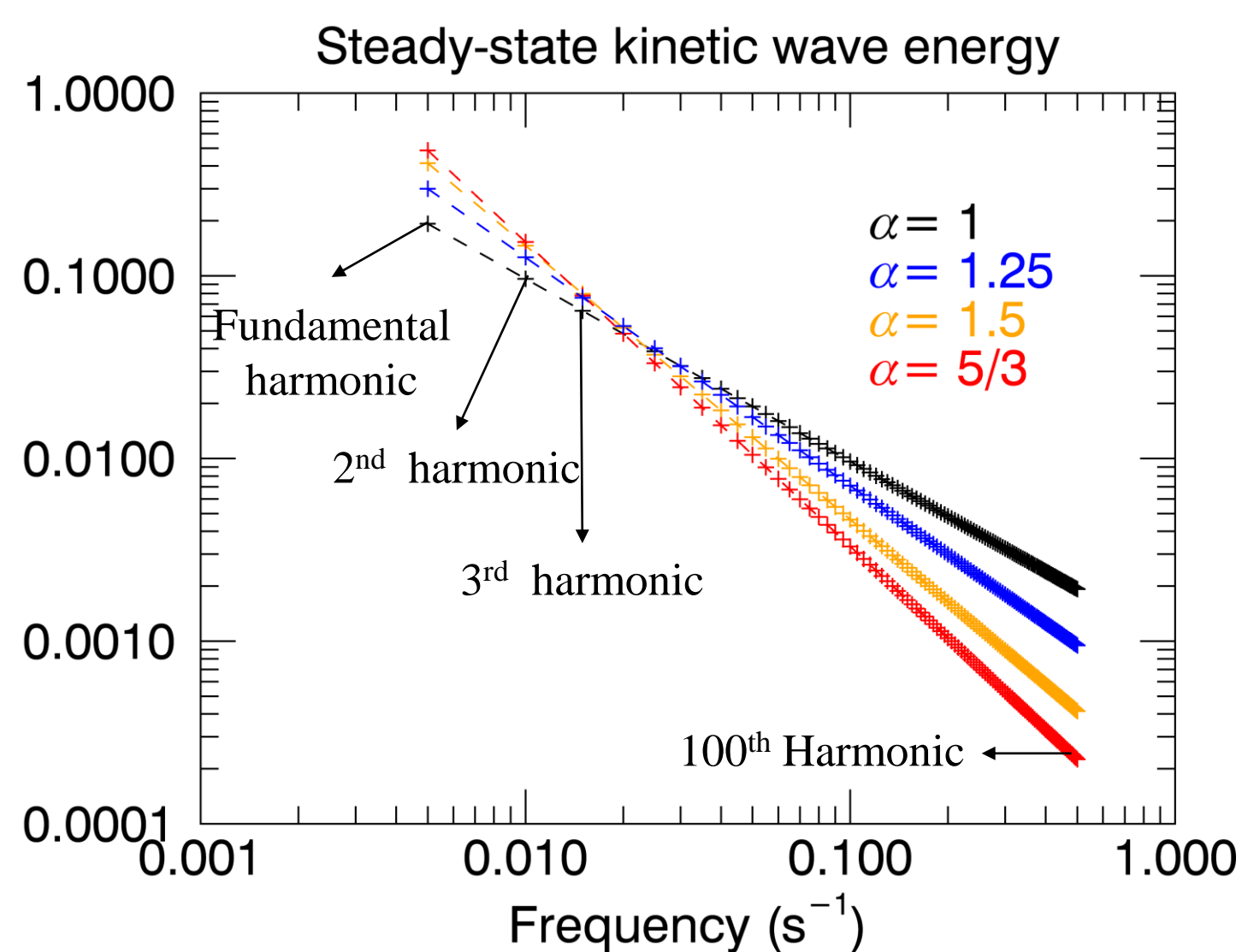


Fig. 2. Normalised power spectrum of the steady-state kinetic wave energy of the resonant field line.

References:

Cargill, P.J., De Moortel, I. and Kiddie, G., 2016. *The Astrophysical Journal*, 823(1), p.31.
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McIntosh, S.W., De Pontieu, B., Carlsson, M., Hansteen, V., Boerner, P. and Goossens, M., 2011. *Nature*, 475(7357), p.477.
Prokopyszyn, A.P.K., Hood, A.W. and De Moortel, I., 2019. *Astronomy & Astrophysics*.

Damping rate

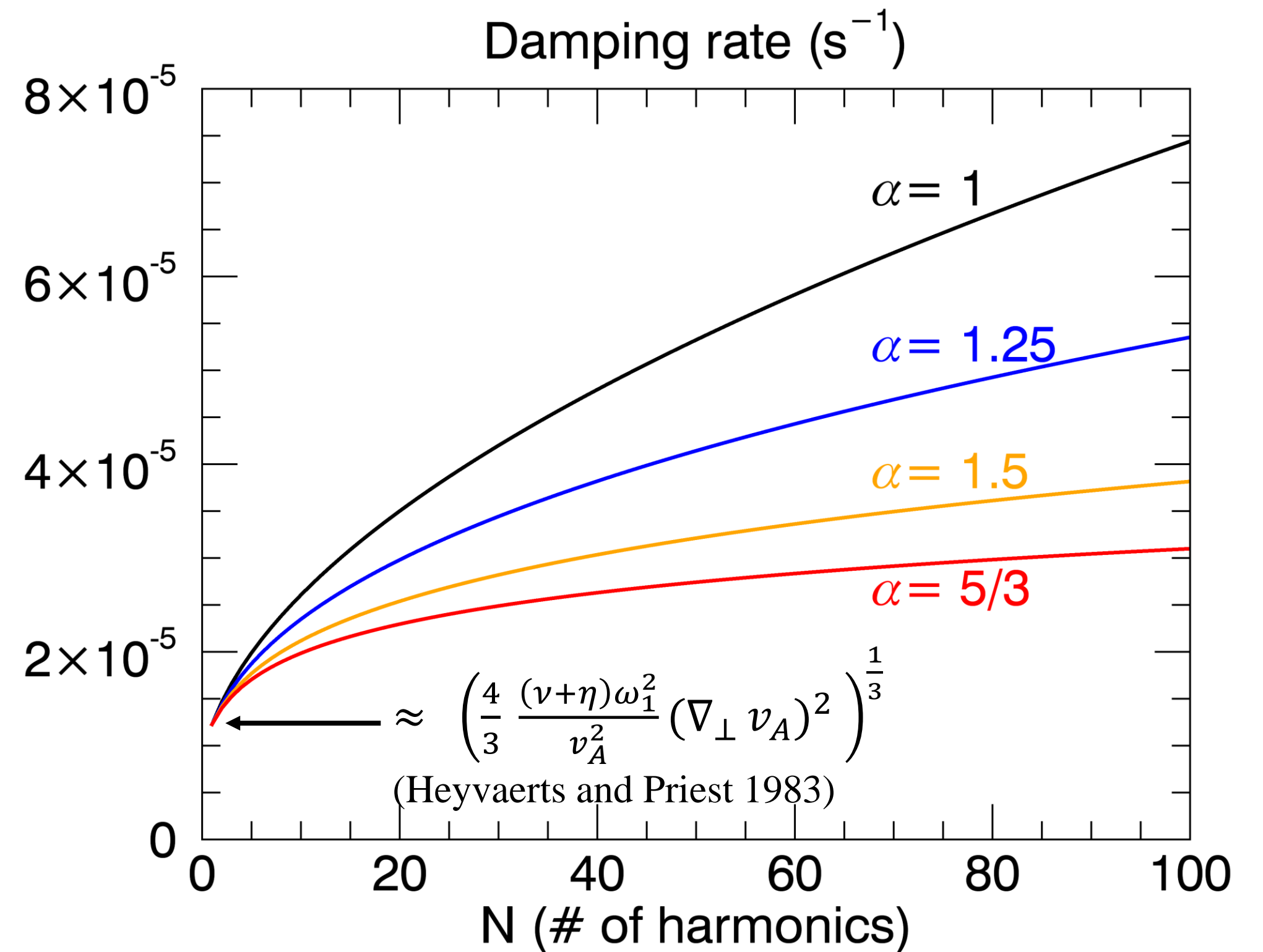


Fig. 3. Damping rate, γ , of the resonant field line as a function of the number of excited harmonics.

- Use $\nu + \eta = 1 \text{ m}^2 \text{ s}^{-1}$, $v_A = 1 \text{ Mm s}^{-1}$, $\nabla_{\perp} v_A = 1 \text{ s}^{-1}$ and $L = 100 \text{ Mm} \Rightarrow$ fundamental frequency of $f_1 = 5 \times 10^{-3} \text{ Hz}$.
- Exciting higher harmonics gives a larger damping rate
- Figures 2 and 3 are independent of the phases ϕ_n .

Upper bound?

- We find that:
 - Leakage of waves through the transition region reduces the damping rate.
 - The damping rate increases with time but converges to a maximum as the system approaches steady-state.
 - The damping rate is largest at resonance.
- Cargill et al. (2016) showed that the thermodynamic response due to the heating reduces the density gradients, which reduces γ .
- Prokopyszyn et al. (2019) showed that nonlinearities have a negligible impact on γ for $u/v_A < 0.1$.
- Relaxing the 2.5D assumption and linear approximation could lead to turbulence which will increase γ , however, we focus here on laminar phase mixing.

\Rightarrow Fig. 3 is an upper bound for γ .

Conclusion

- Using a coronal heating rate, $H_c \approx 10^{-5} \text{ W m}^{-3}$ and observed amplitudes in the quiet sun $\approx 20 \text{ km s}^{-1}$ (McIntosh et al. 2011)

\Rightarrow we require $\gamma \approx 10^{-1} \text{ s}^{-1}$.

- Figure 3 shows γ is too small by 3 orders of magnitude.
- **If our estimate is an upper bound \Rightarrow laminar phase mixing is not a viable standalone heating mechanism in coronal loops.**
- Laminar phase mixing may be significant in other setups, for example near null points where the Alfvén speed is small, and the cross-field viscosity is stronger.