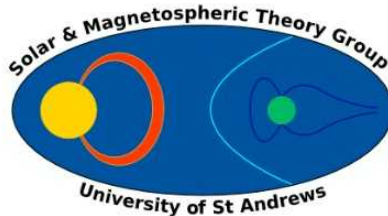


Mode coupling at the transition region and the validity of line-tied boundary conditions

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University of
St Andrews



Science & Technology
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Aims

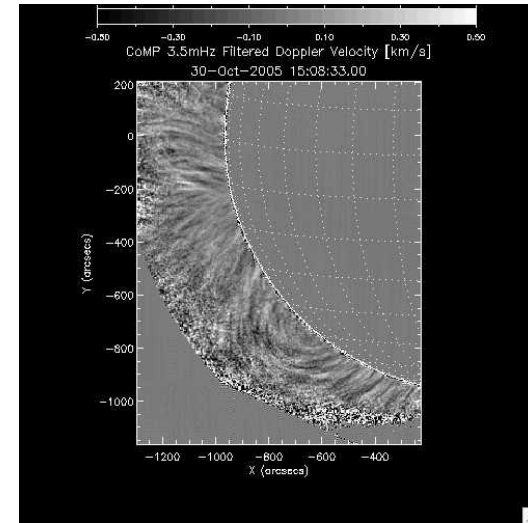
- Show why Fast / Alfvén waves couple at the TR
- Show that polarisation of the waves changes upon reflection
- Test the validity of line-tied BCs

Structure

- **Background**
- Model 1:
 - Line-tied, pulse
- Model 2:
 - Line-tied, normal mode
- Model 3:
 - Chromosphere, normal mode
- Summary and conclusions

Why study MHD waves?

- Ubiquitous
- Coronal heating
- Coronal seismology

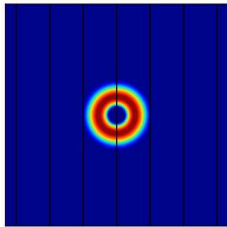


Tomczyk et al. (2007)

Fast vs. Alfvén waves

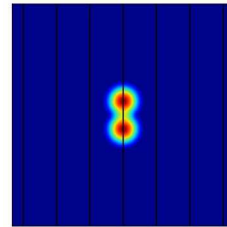
Fast waves:

- Propagate isotropically
- $\frac{\omega}{v_A} = \pm \sqrt{k_x^2 + k_y^2 + k_z^2}$



Alfvén waves:

- Propagate parallel to \mathbf{B}_0
- $\frac{\omega}{v_A} = \pm k_{\parallel}$



Mode conversion

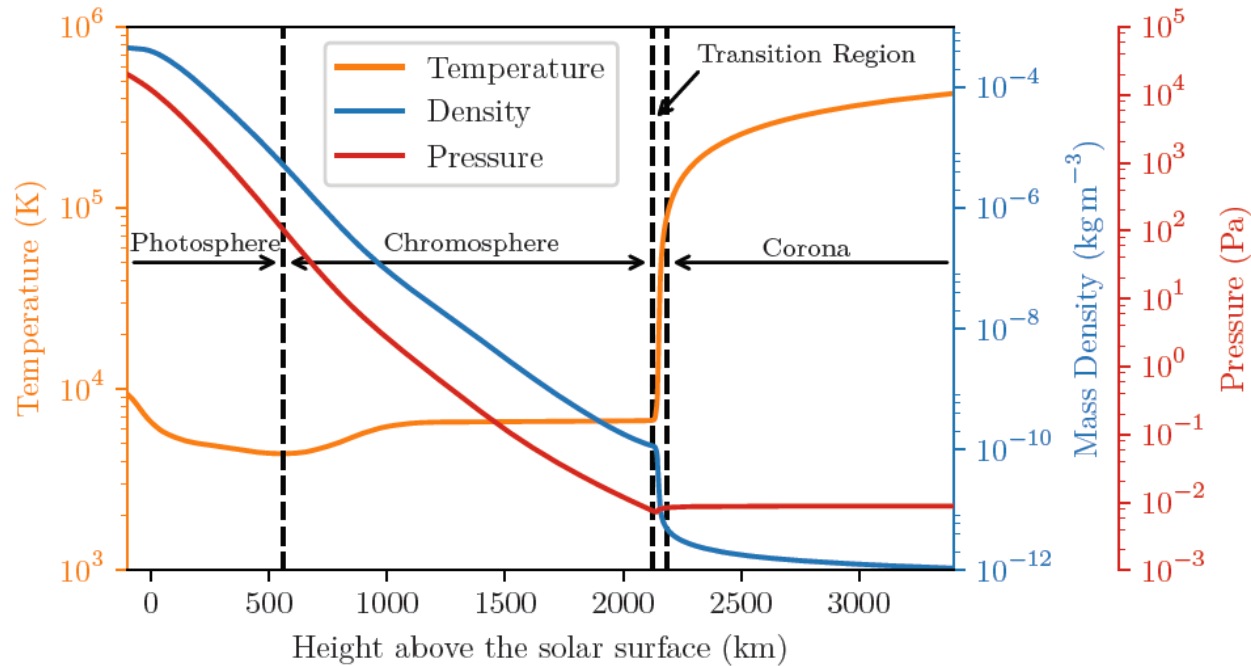
Can occur via:

- Non-linear effects (Verwichte et al., 1999)
- Transition from $\beta < 1$ to $\beta > 1$ plasma (McLaughlin & Hood, 2006)
- Gradients in $v_A \rightarrow$ resonant absorption (Ionson, 1982)

Mode conversion at the TR

- Studied analytically in Halberstadt & Goedbloed (1993, 1994, 1995)
- Numerical approach used in Arregui et al. (2003)
- Cally & Hansen (2011, 2012) suggest that mode conversion from fast waves to Alfvén waves at the transition region enables sufficient energy flux to enter the corona

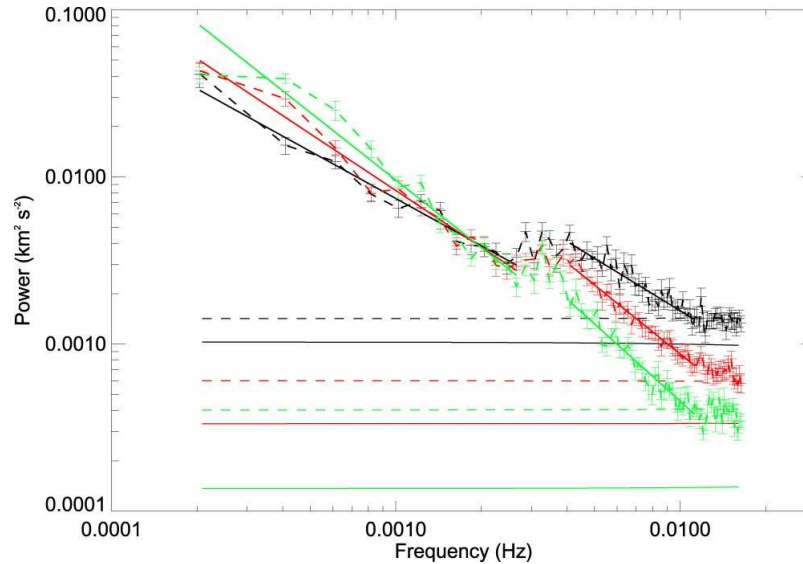
Line-tied ($u=0$) boundary conditions



Vernazza et al. (1981) and Williams (2018)

Normal mode

- $f(\mathbf{r}, t) = f_0(\mathbf{r})\exp(i\omega t)$



Morton et al. (2016)

Model and Equations

- Background quantities:

$$\rho = \rho_0$$

$$\mathbf{B}_0 = B_0 \hat{\mathbf{B}}_0$$

- Perturbations:

$$\mathbf{u} = u_x \hat{\mathbf{x}} + u_{\perp} \hat{\perp}$$

$$\mathbf{b} = b_x \hat{\mathbf{x}} + b_{\perp} \hat{\perp} + b_{\parallel} \hat{\mathbf{B}}_0$$

- Unit vectors:

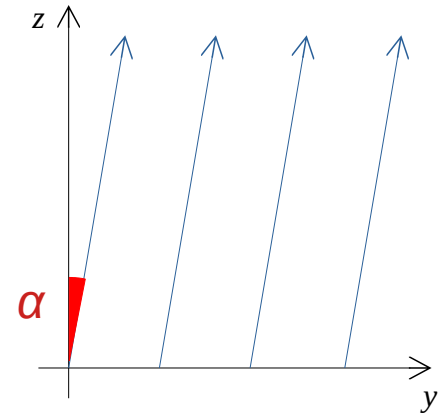
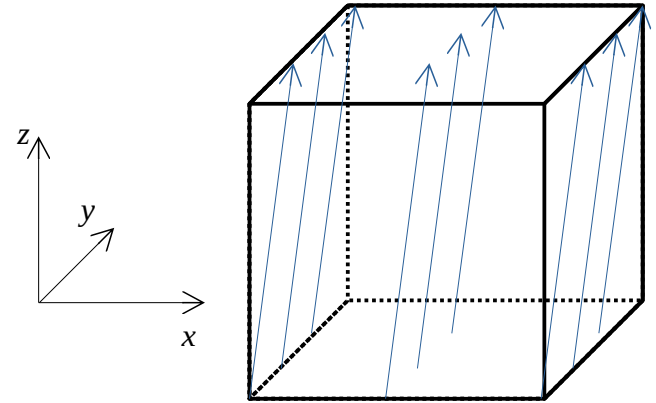
$$\hat{\perp} = \cos(\alpha) \hat{\mathbf{y}} - \sin(\alpha) \hat{\mathbf{z}}$$

$$\hat{\mathbf{B}}_0 = \sin(\alpha) \hat{\mathbf{y}} + \cos(\alpha) \hat{\mathbf{z}}$$

- Equations:

$$\rho_0 \frac{\partial \mathbf{u}}{\partial t} = \mathbf{j} \times \mathbf{B}_0$$

$$\frac{\partial \mathbf{b}}{\partial t} = \nabla \times (\mathbf{u} \times \mathbf{B}_0)$$



Structure

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Numerical scheme

- Leapfrog algorithm
- Based on Zalesak (1979)
- Finite-difference
- Staggered grid
- Second-order accurate

Initial / boundary conditions

- Assume that

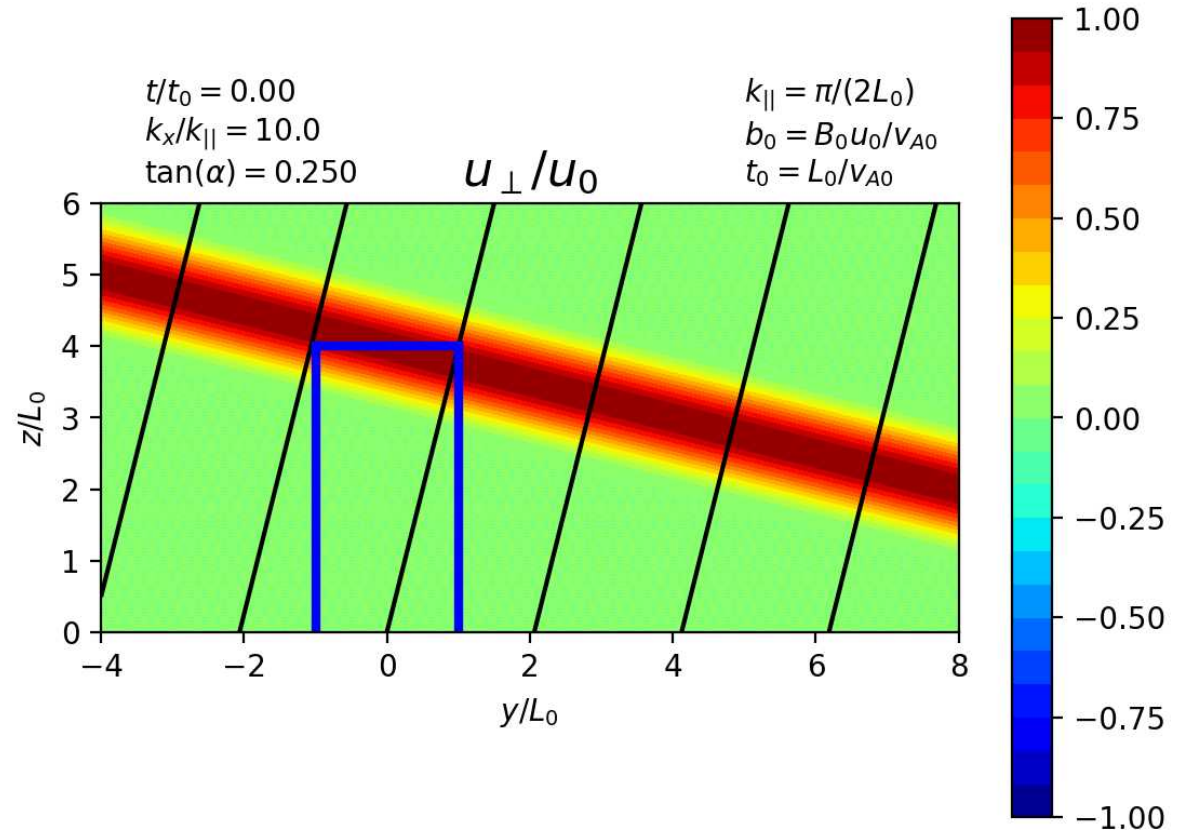
$$u_x, b_x \propto \sin(k_x x)$$

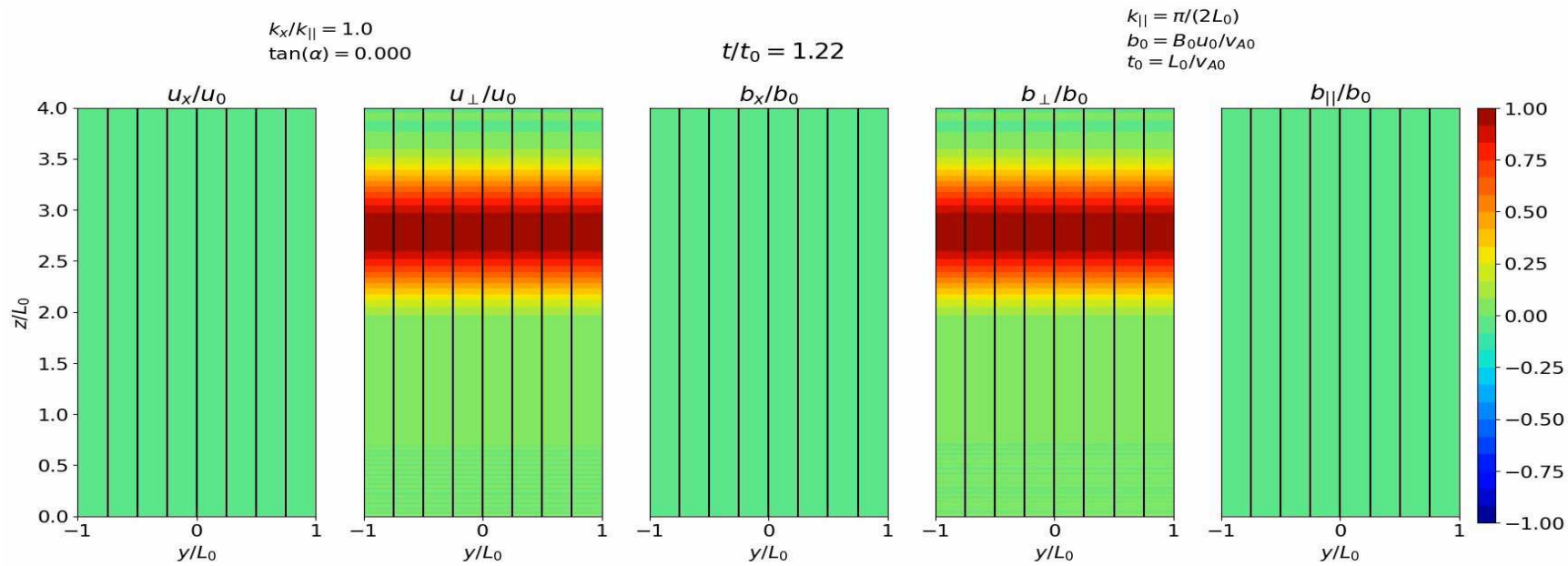
$$u_\perp, u_\parallel, b_\parallel \propto \cos(k_x x)$$

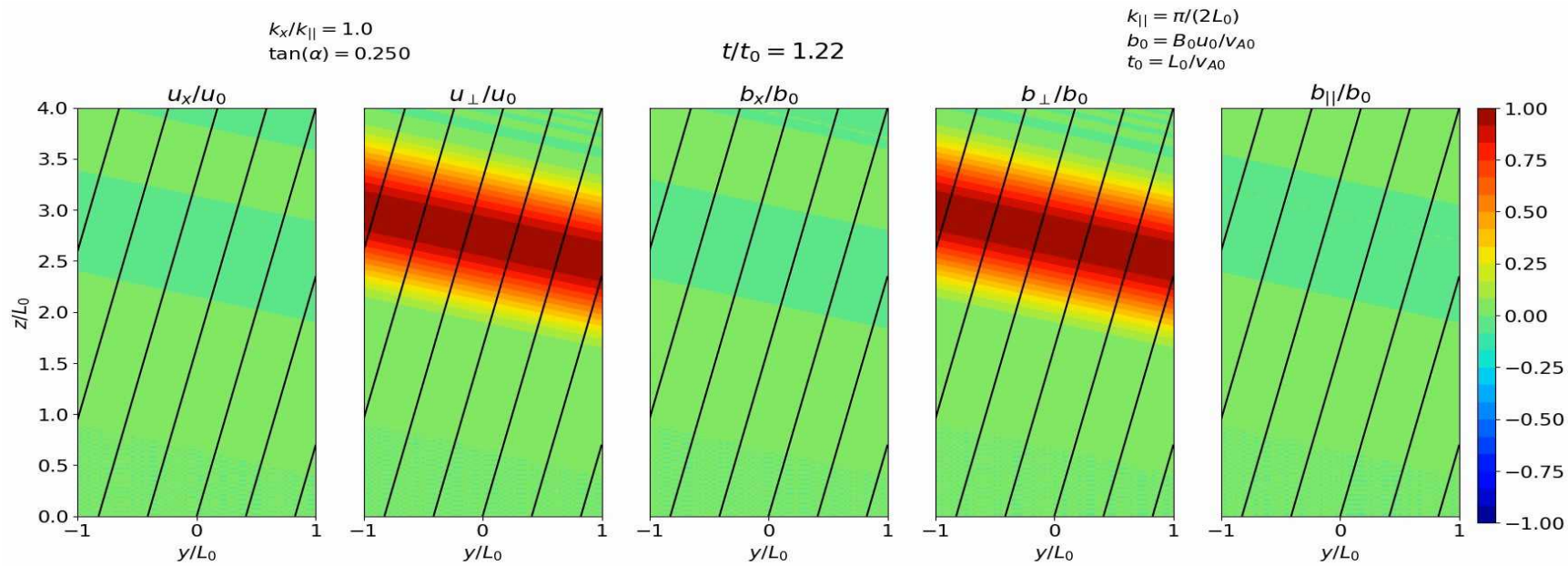
- Initial conditions:

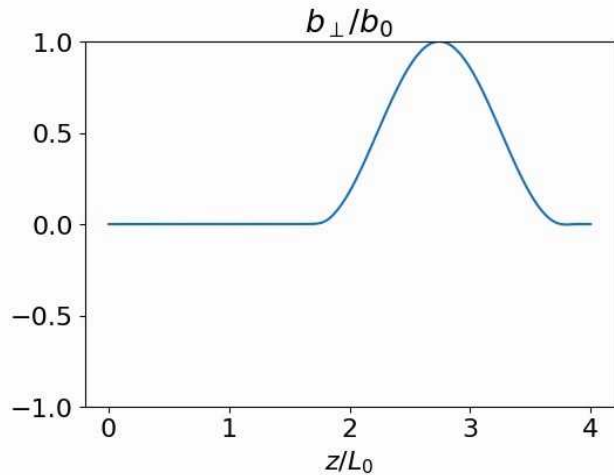
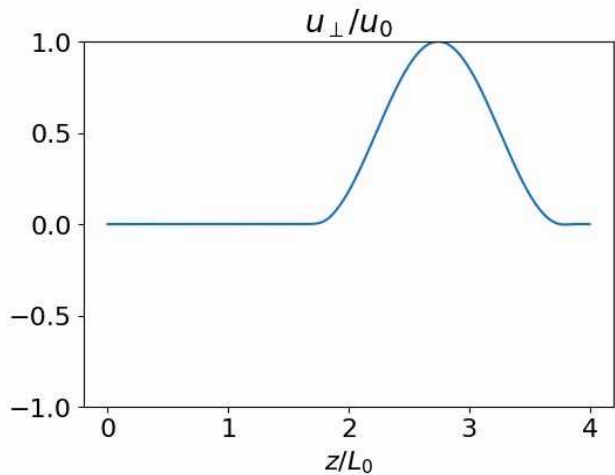
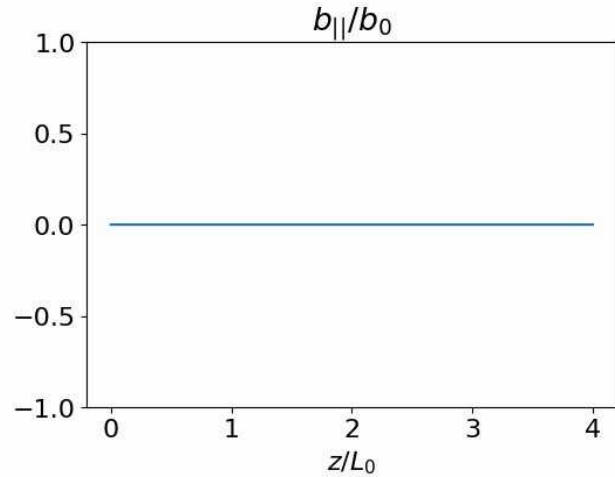
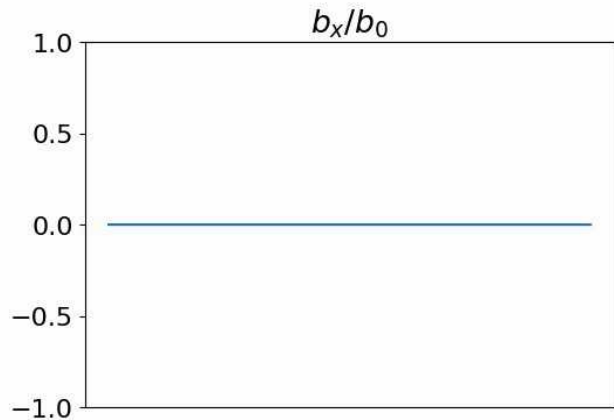
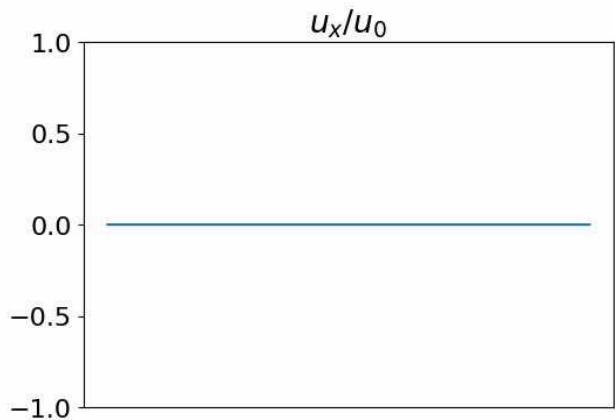
$$\frac{u_\perp}{u_0} = \frac{b_\perp}{b_0} = \begin{cases} \cos^2 \theta & \text{if } |\theta| \leq \pi/2 \\ 0 & \text{if } |\theta| > \pi/2 \end{cases}$$

$$\theta = k_\parallel (y \sin \alpha + (z + 4L_0) \cos \alpha + v_A t)$$









$$t/t_0 = 1.22$$

$$y = 0.0$$

$$k_x/k_{||} = 1.0$$

$$\tan(\alpha) = 0.250$$

$$k_{||} = \pi/(2L_0)$$

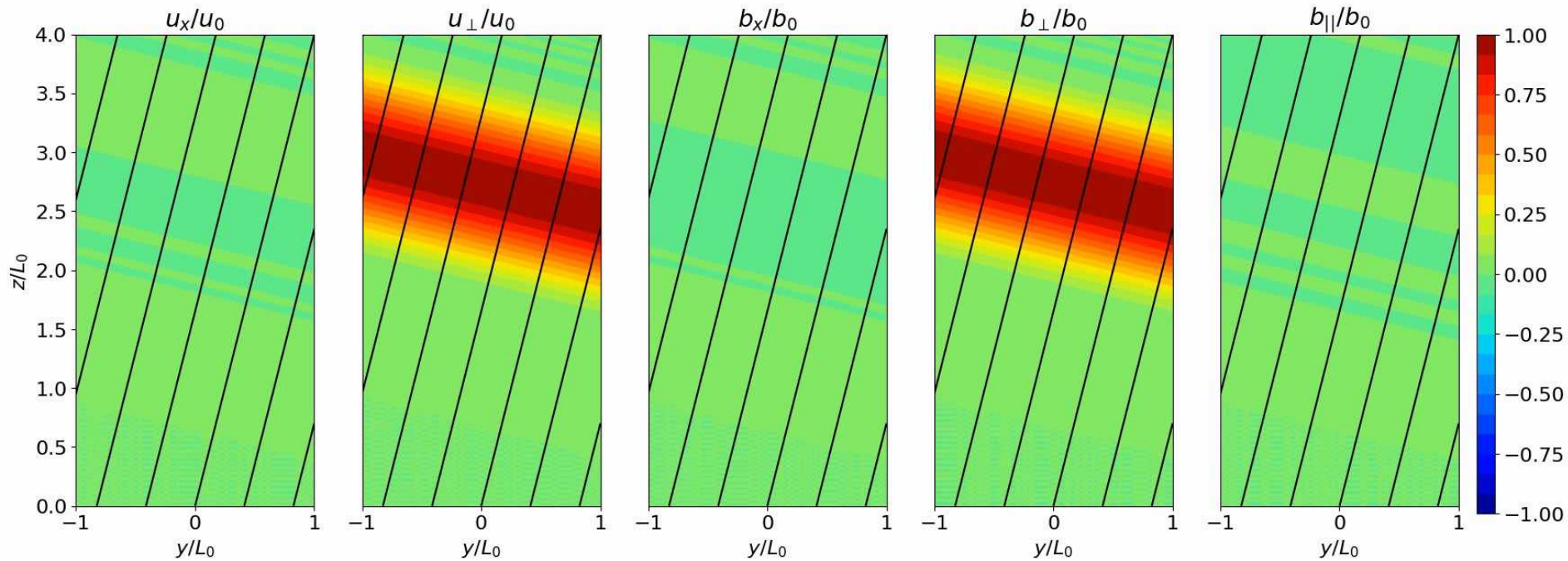
$$b_0 = B_0 u_0 / v_{A0}$$

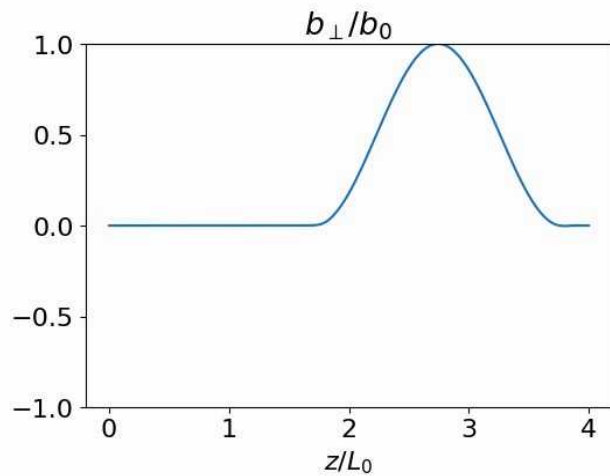
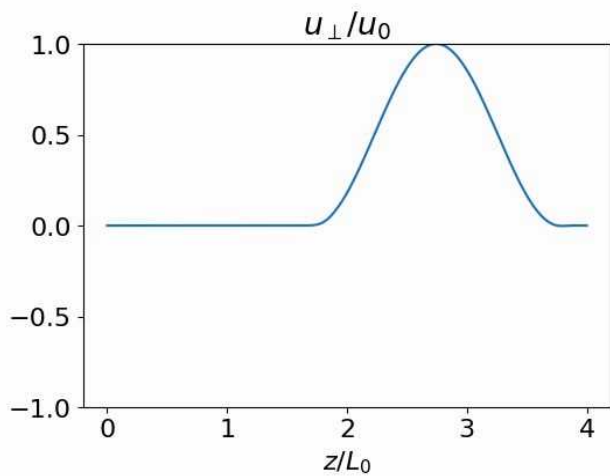
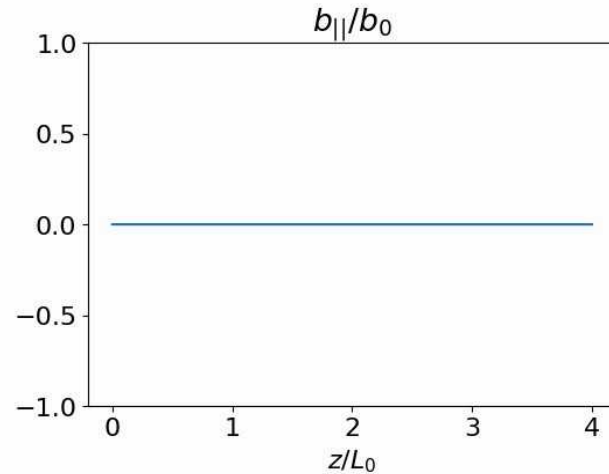
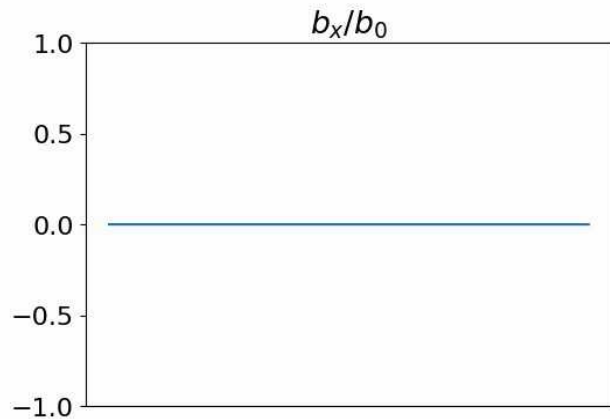
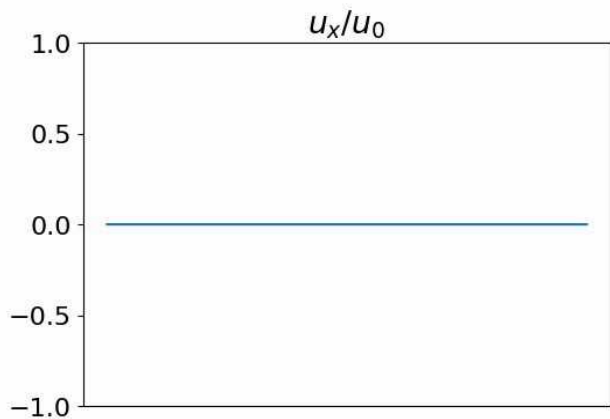
$$t_0 = L_0 / v_{A0}$$

$k_x/k_{\parallel} = 10.0$
 $\tan(\alpha) = 0.250$

$t/t_0 = 1.22$

$k_{\parallel} = \pi/(2L_0)$
 $b_0 = B_0 u_0 / v_{A0}$
 $t_0 = L_0 / v_{A0}$





$$t/t_0 = 1.22$$

$$y = 0.0$$

$$k_x/k_{||} = 10.0$$

$$\tan(\alpha) = 0.250$$

$$k_{||} = \pi/(2L_0)$$

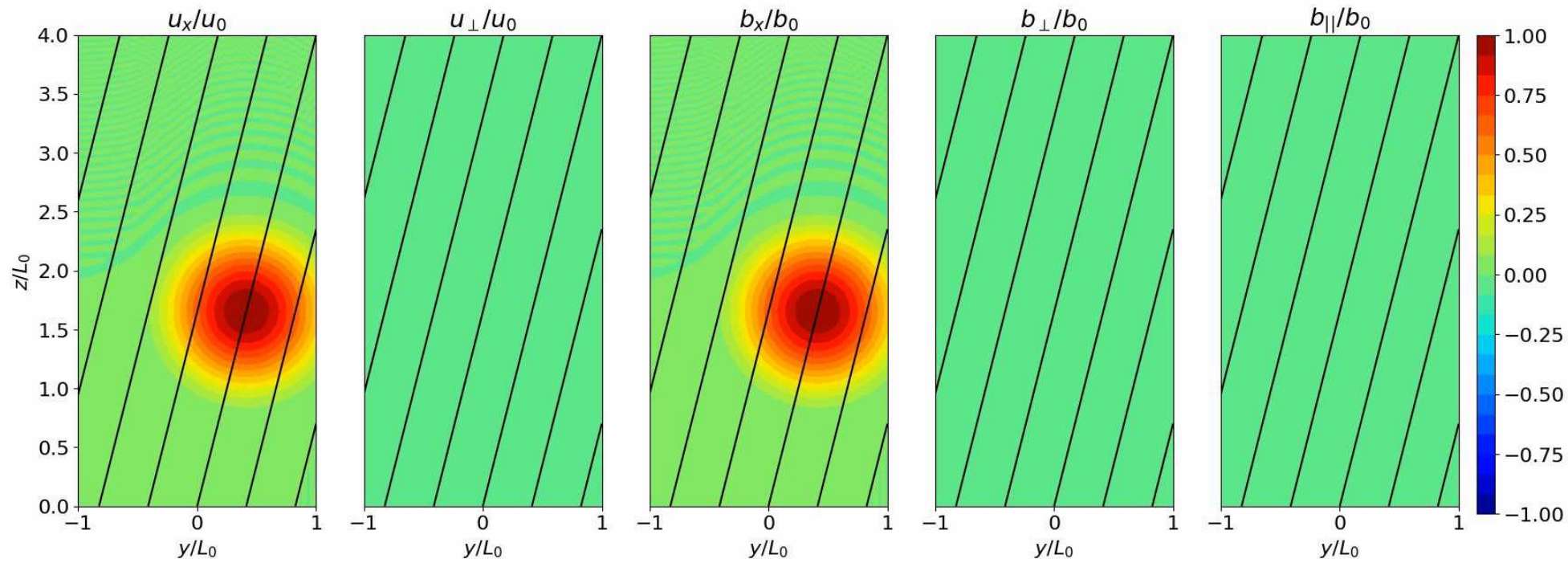
$$b_0 = B_0 u_0 / v_{A0}$$

$$t_0 = L_0 / v_{A0}$$

$k_x/k_{||} = 0.0$
 $\tan(\alpha) = 0.250$

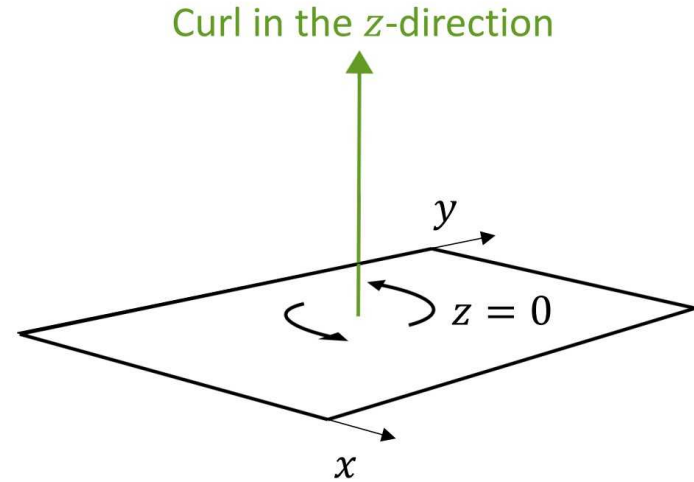
$t/t_0 = 2.42$

$k_{||} = \pi/(2L_0)$
 $b_0 = B_0 u_0 / v_{A0}$
 $t_0 = L_0 / v_{A0}$



Why does the coupling occur?

- $\frac{\partial b_z}{\partial t} = \hat{z} \cdot \nabla \times (\mathbf{u} \times \mathbf{B}_0)$
- $b_z = 0$ at $z = 0$
- $b_z = \cos(\alpha)b_{\parallel} - \sin(\alpha)b_{\perp}$
- $b_{\parallel} = \tan(\alpha)b_{\perp} \Rightarrow$ Fast waves



Summary

At the solar surface:

- Alfvén waves couple to fast waves
- Change polarisation
- If

$$k_x^2 > k_{\parallel}^2 - k_y^2$$

then evanescent boundary layers form

Structure

- Background
- Model 1:
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Normal mode solution

- Assume

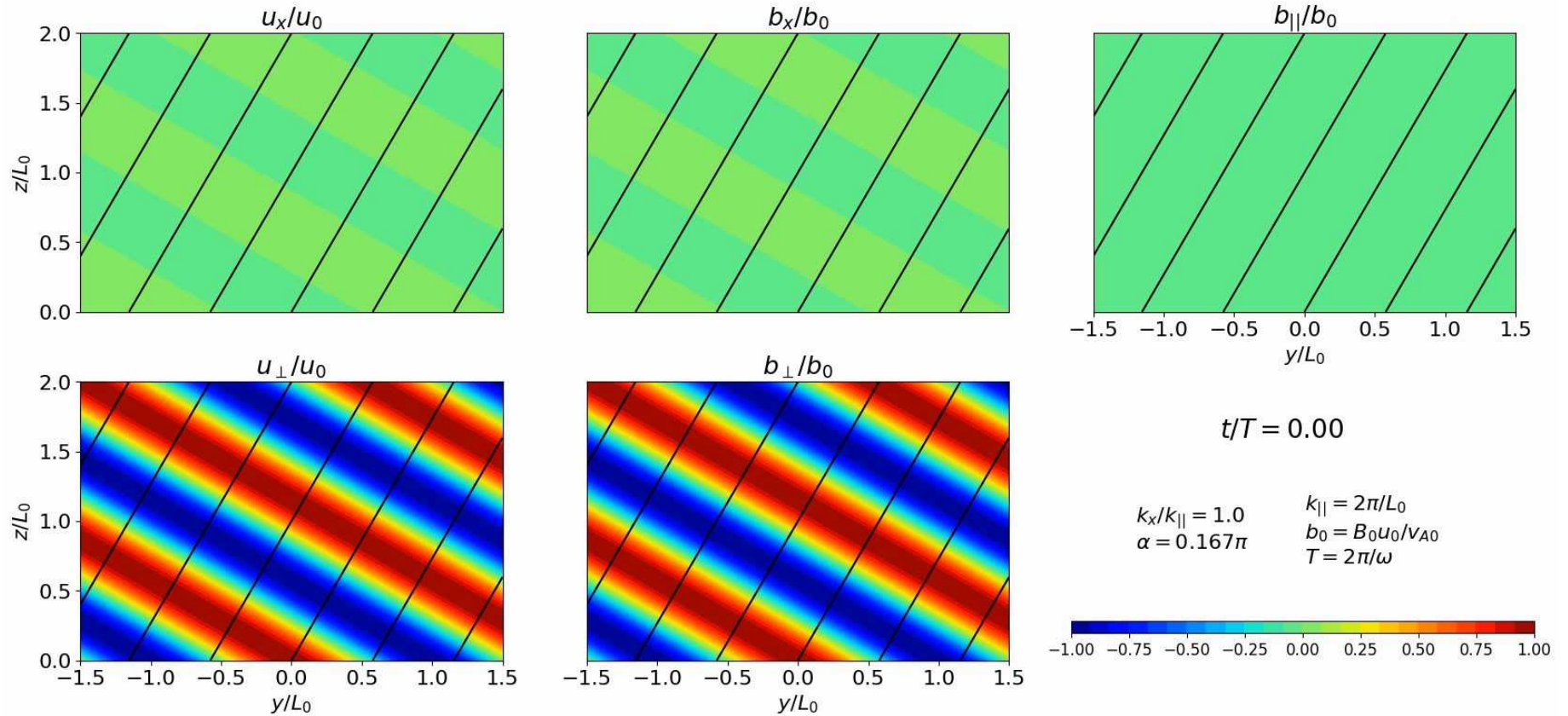
$$u_x, u_\perp, b_x, b_\perp, b_\parallel \propto \exp[i(k_x x + k_y y + \omega t)]$$

- Impose incident Alfvén wave

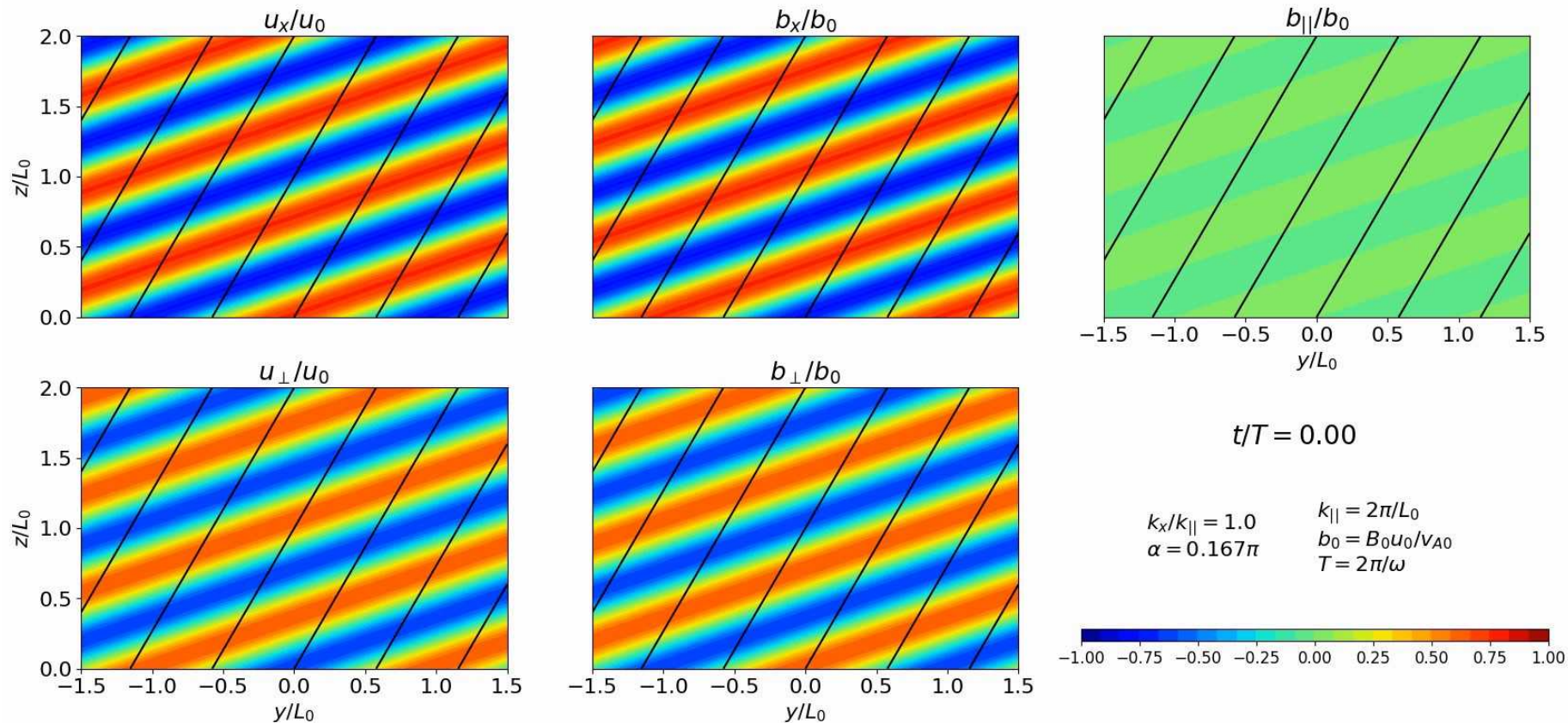
$$\frac{u_\perp}{u_0} = \frac{b_\perp}{b_0} = \exp[i(k_x x + k_\parallel s) + \omega t]$$

- Calculate unique reflected Alfvén and fast wave which ensures $\mathbf{u} = \mathbf{0}$

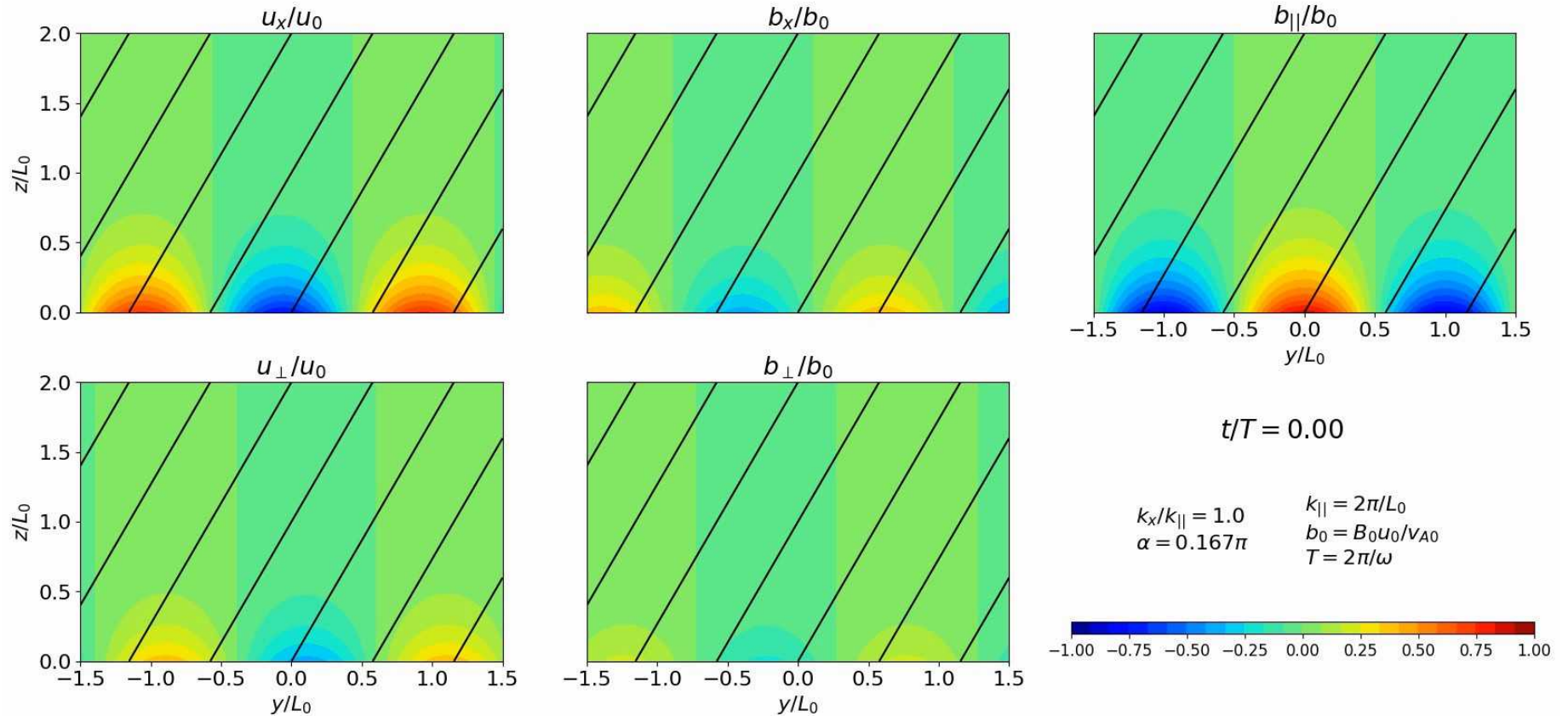
Incident wave



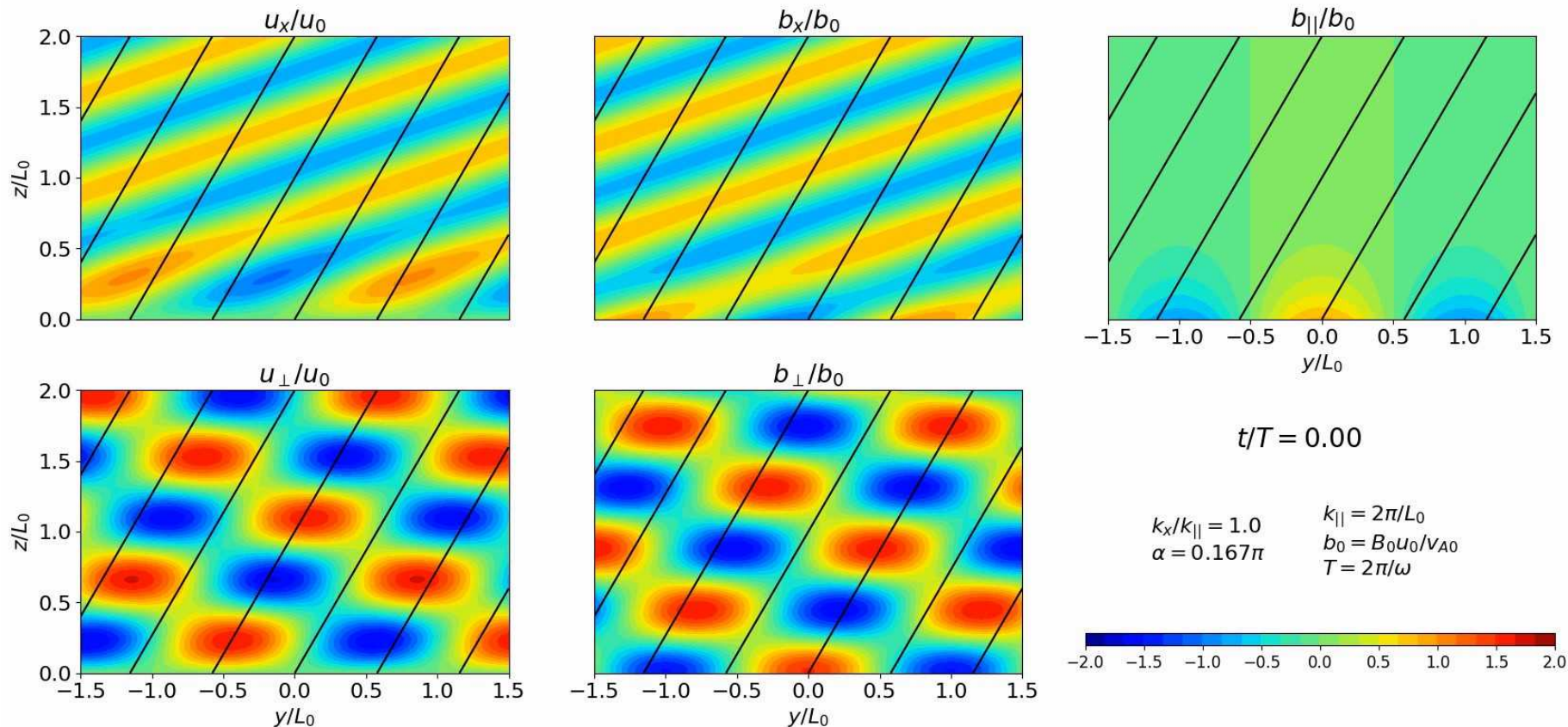
Reflected Alfvén wave



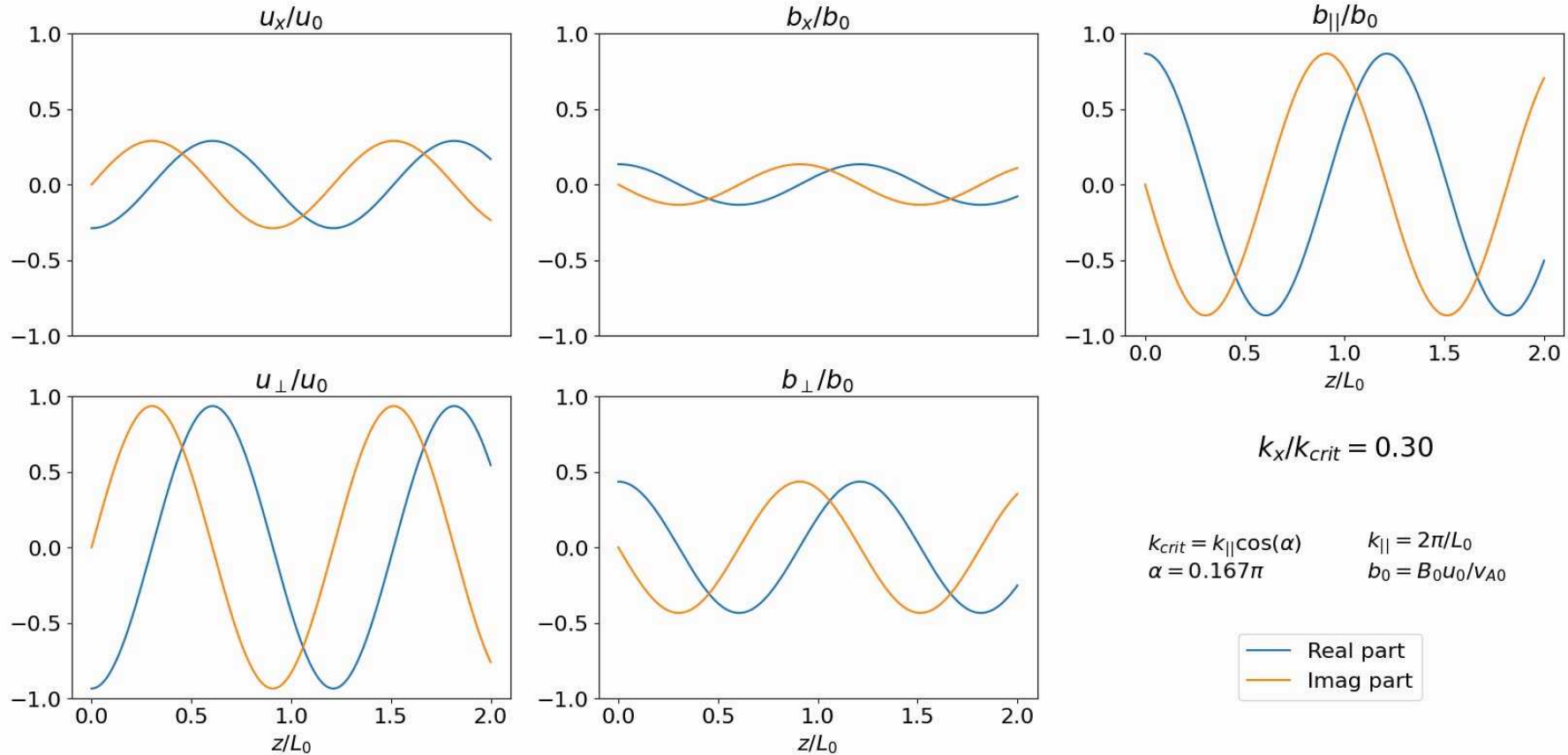
Reflected Fast wave



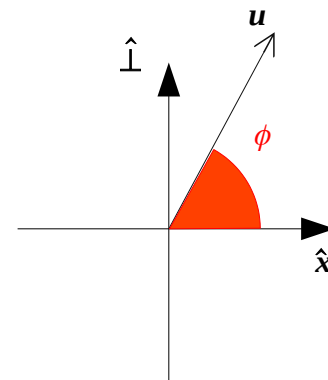
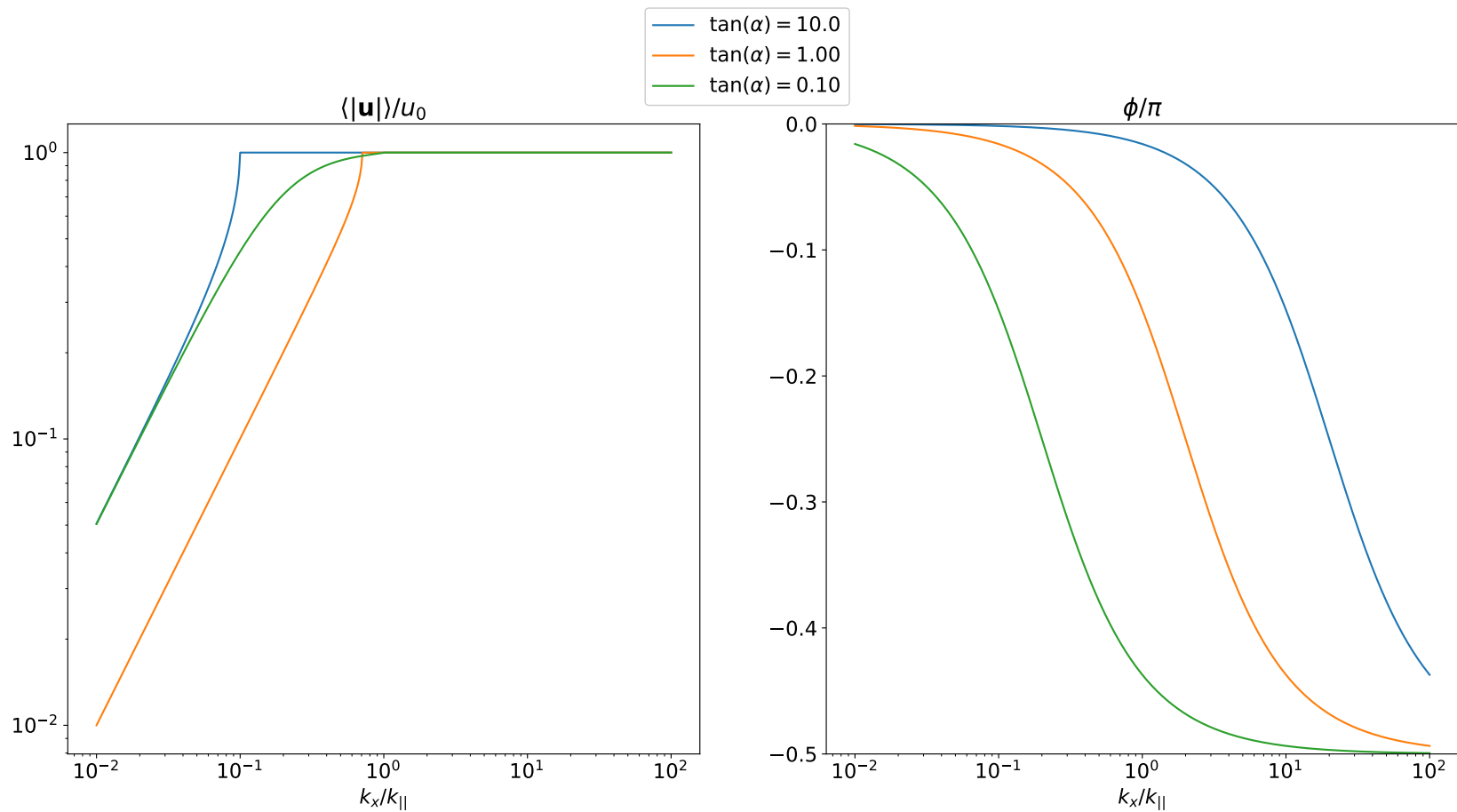
Full solution



Reflected fast wave



Reflected Alfvén wave



Summary

- Fast wave energy $\rightarrow 0$ as $k_x \rightarrow \infty$
- Change in polarisation $\rightarrow 0$ as $k_x \rightarrow \infty$
- \therefore Boundary layers have a minimal impact on resonance absorption

Structure

- Background
- Model 1:
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 - Line-tied, normal mode
- **Model 3:**
 - **Chromosphere, normal mode**
- Summary and conclusions

Model

- Background Alfvén speed:

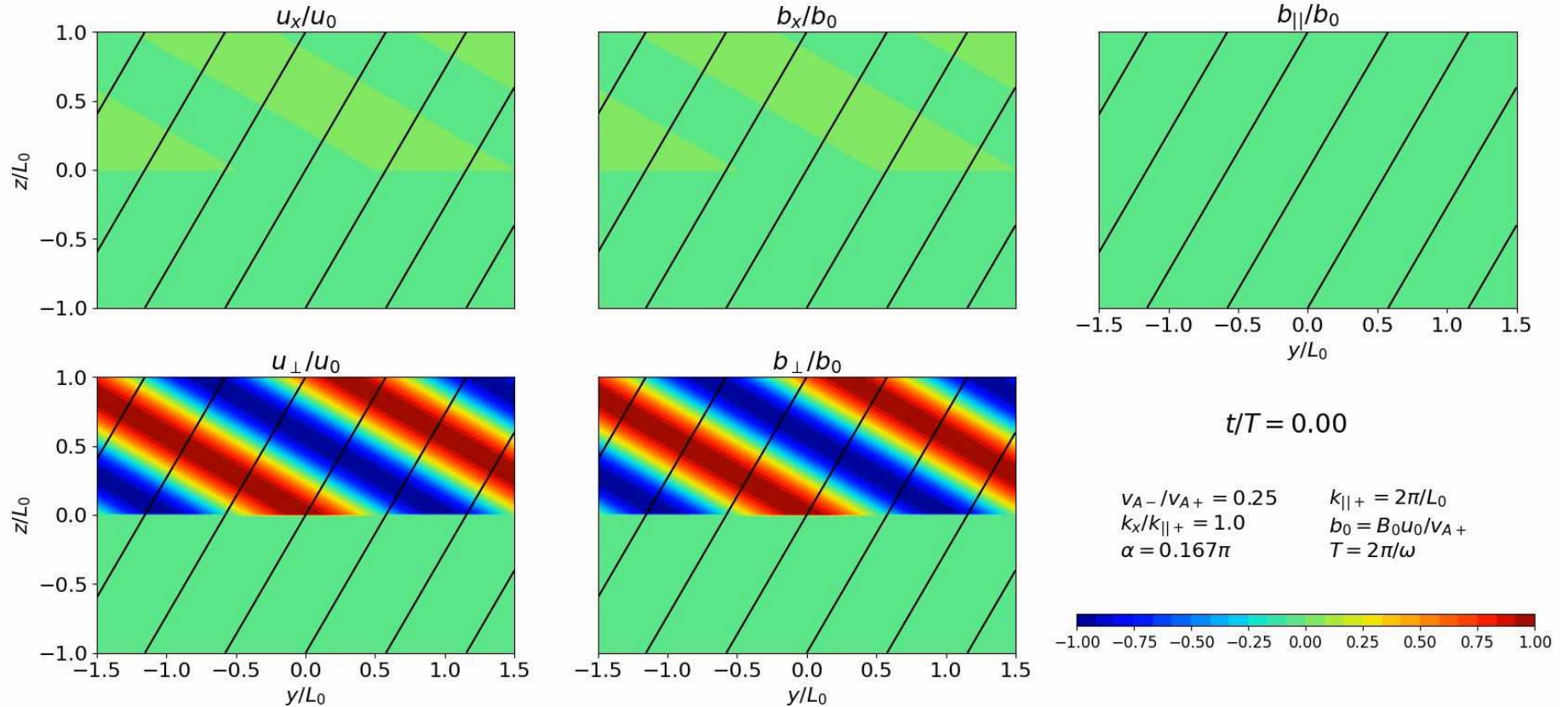
$$v_A = \begin{cases} v_{A+} & \text{if } z \geq 0 \\ v_{A-} & \text{if } z < 0 \end{cases} \quad \begin{array}{l} \leftarrow \text{Corona} \\ \leftarrow \text{Chromosphere} \end{array}$$

- Where:

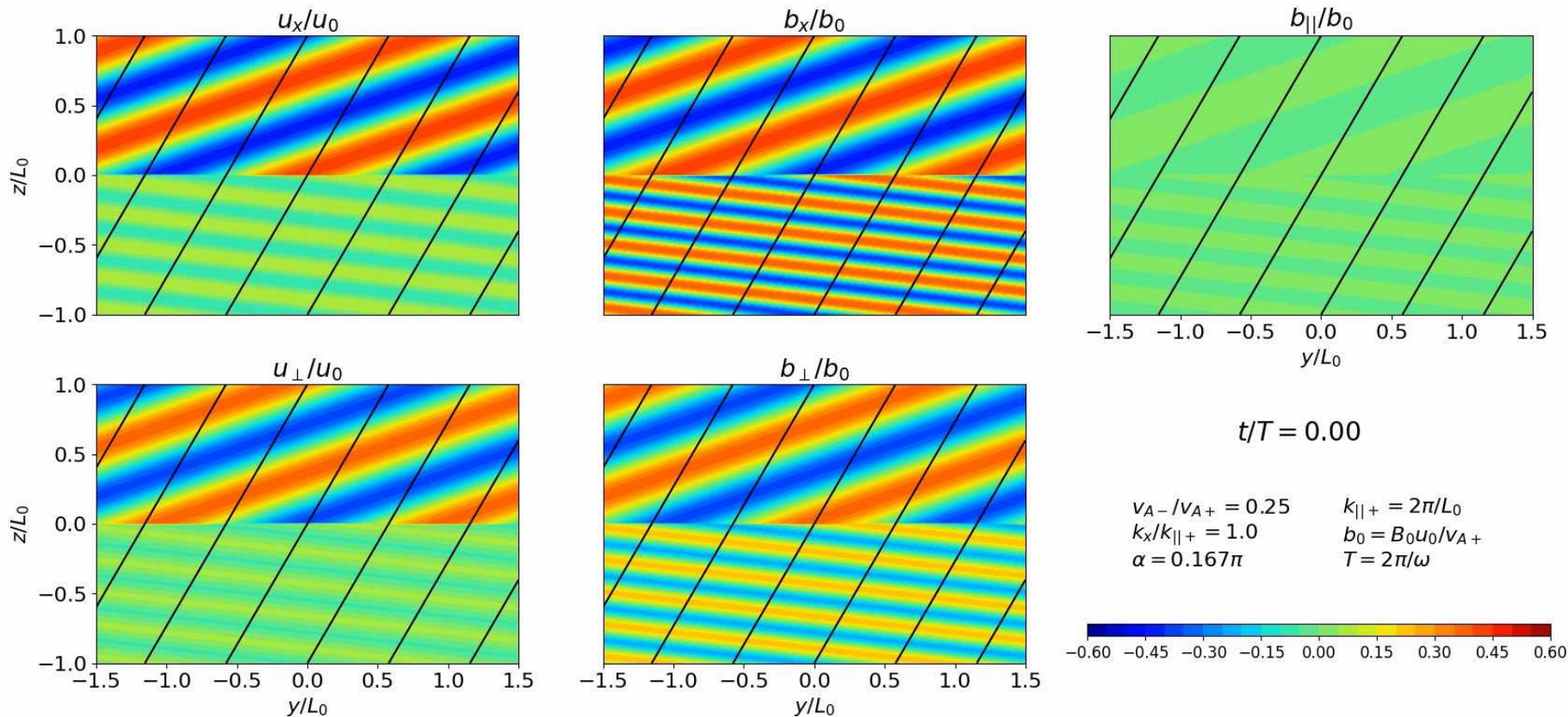
$$v_{A+} \gg v_{A-} \Rightarrow k_{\parallel-} \gg k_{\parallel+}$$

- Impose continuity of \mathbf{u} and \mathbf{b}

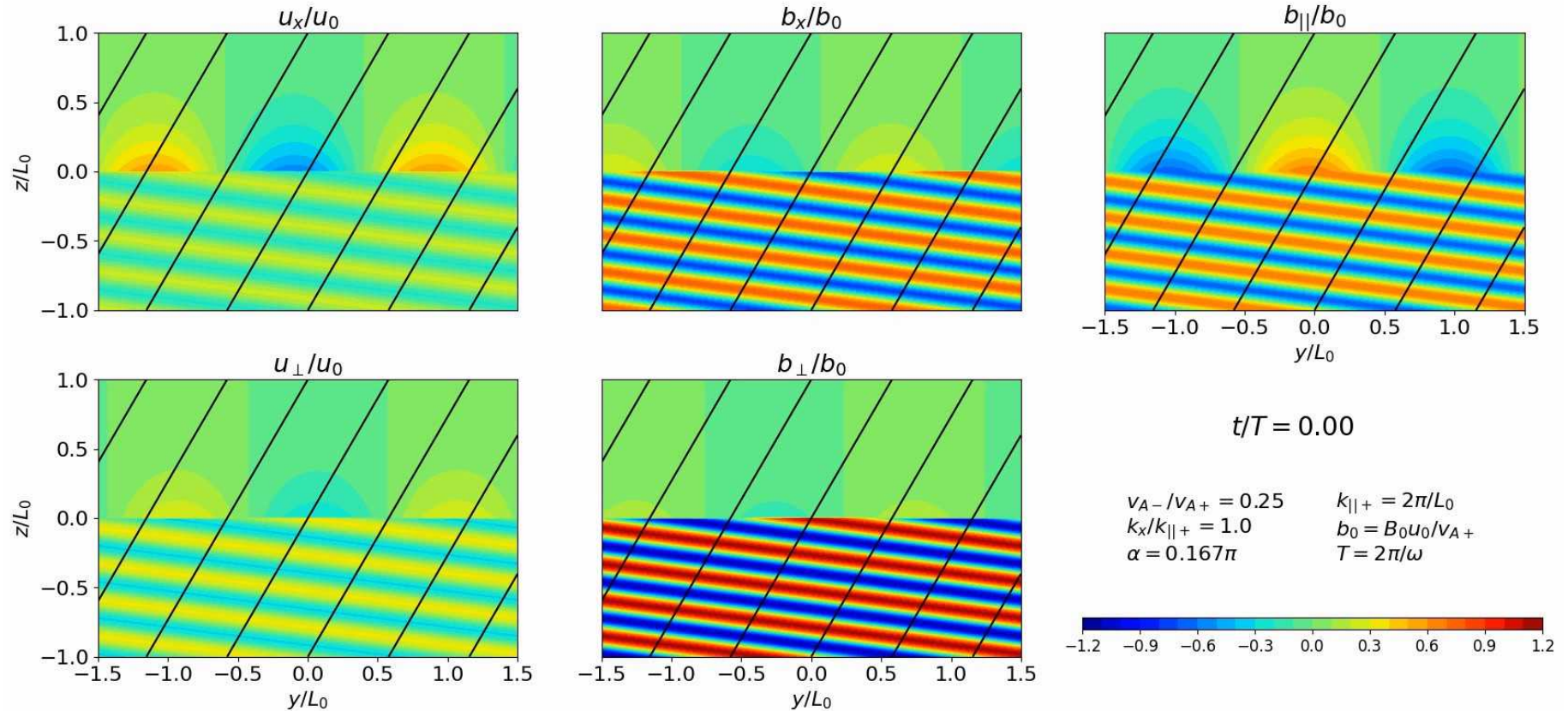
Incident wave



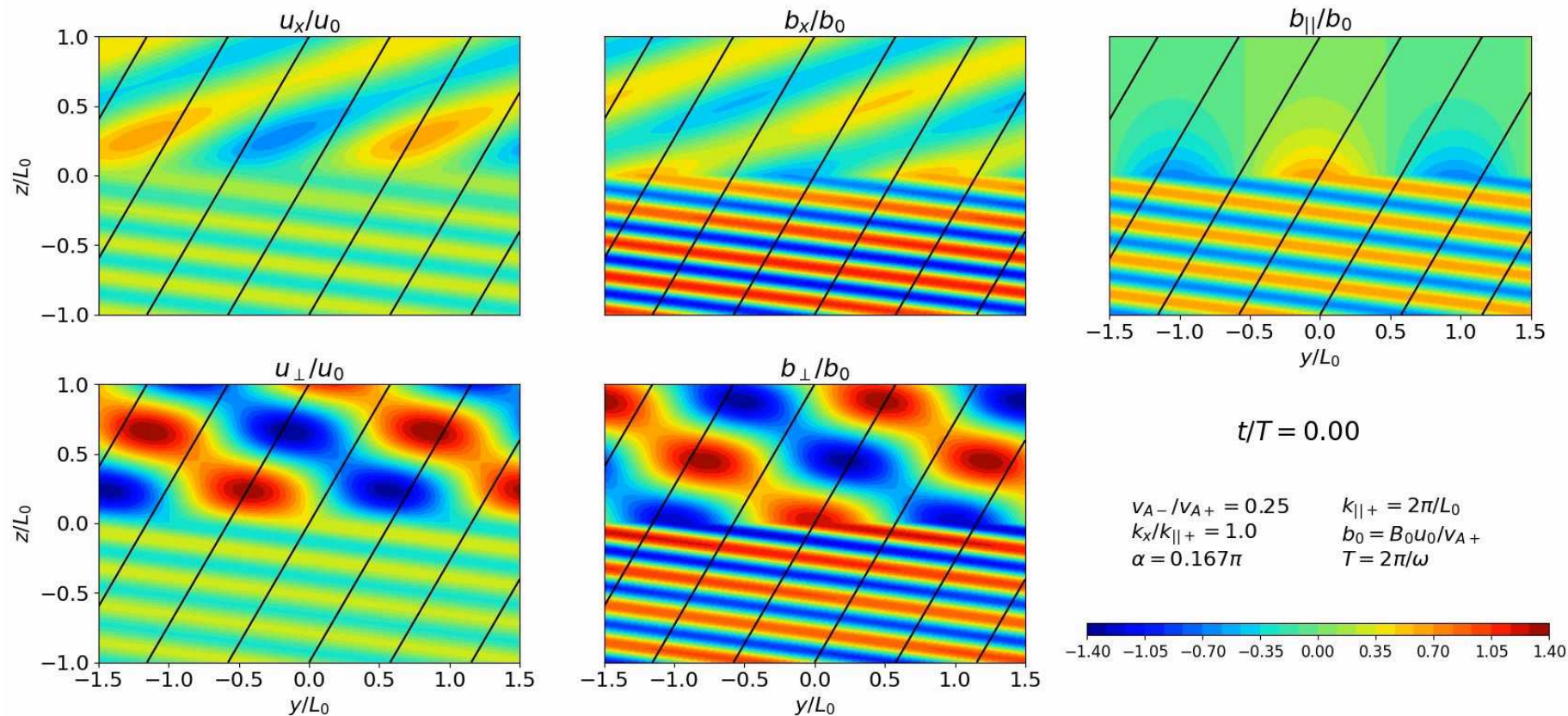
Reflected + Transmitted Alfvén wave



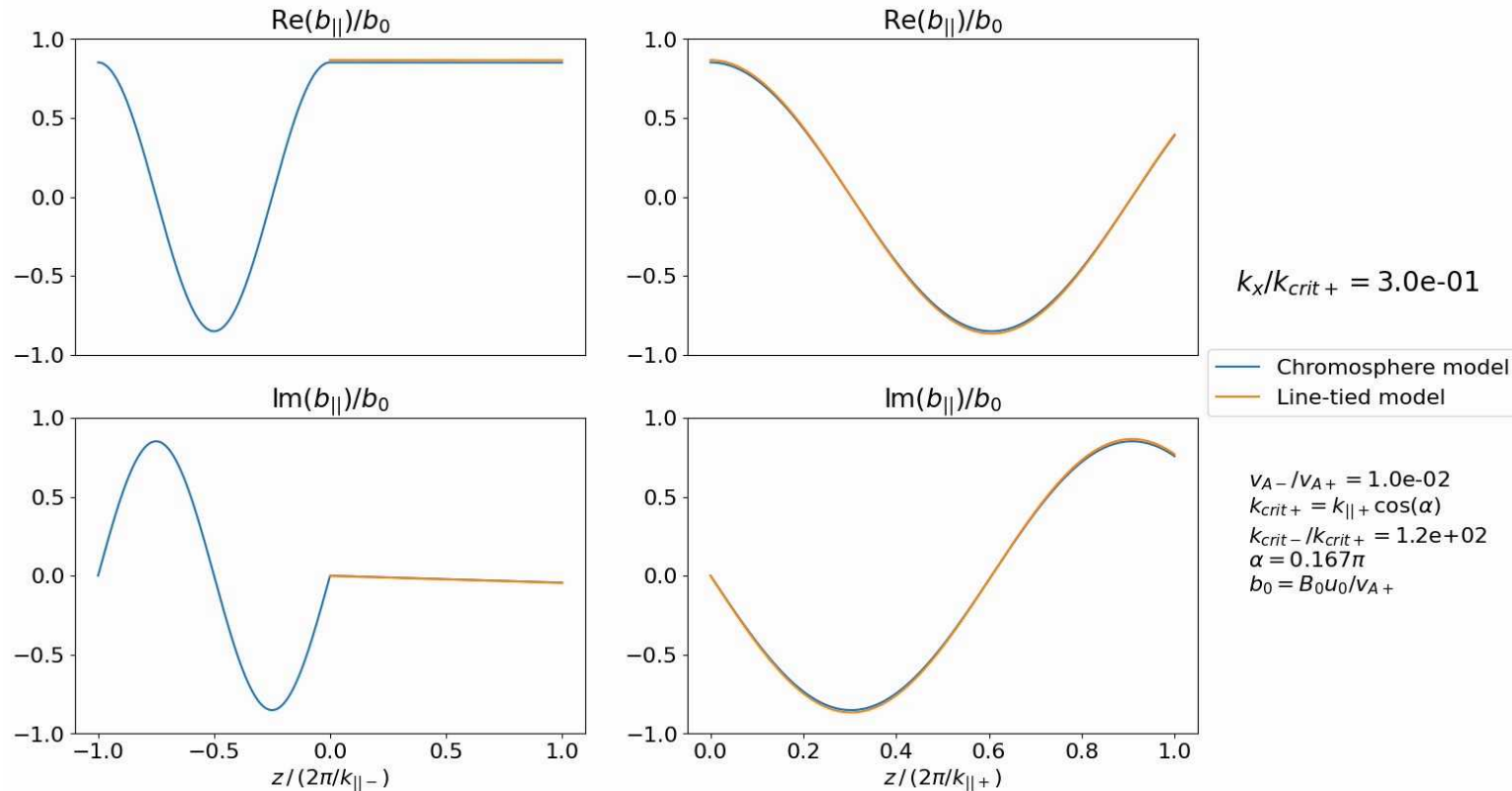
Reflected + Transmitted Fast wave



Full solution

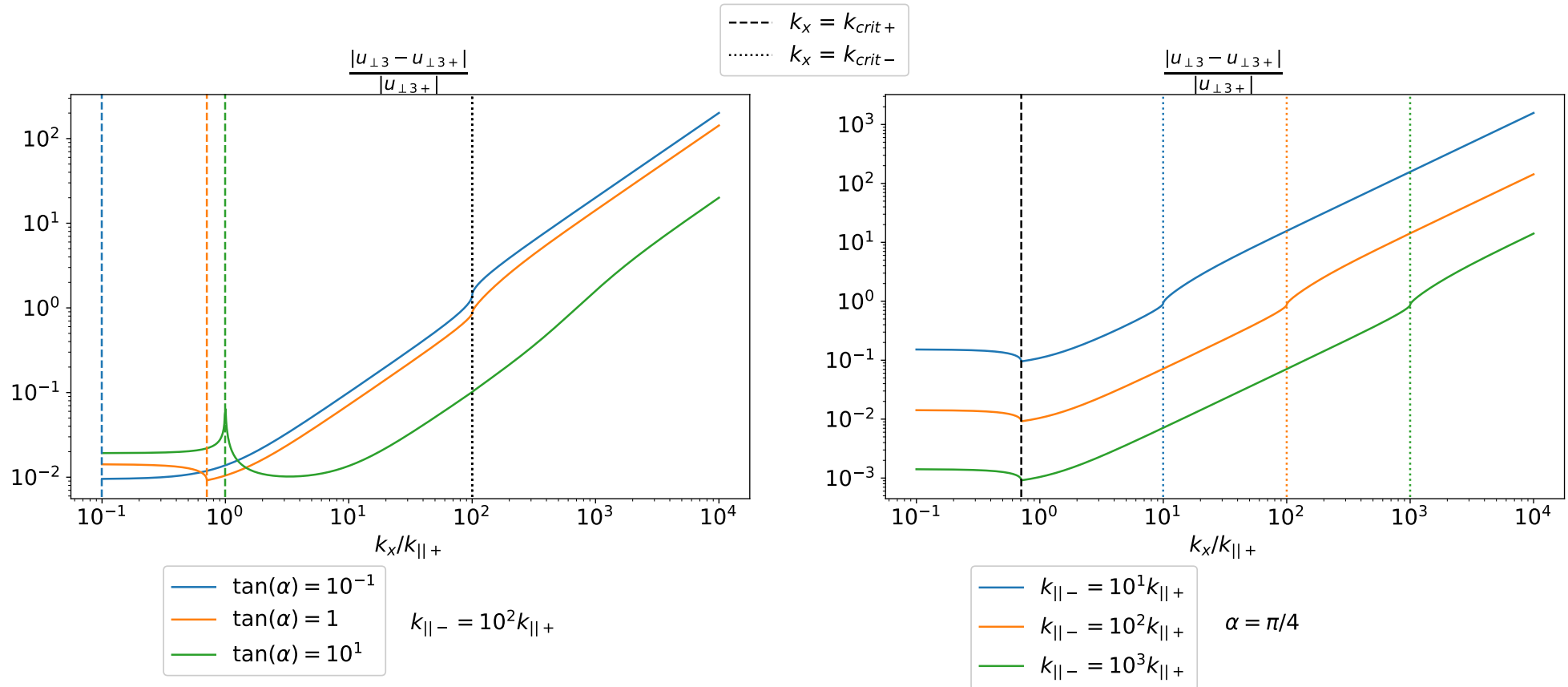


Chromosphere vs. Line-tied model

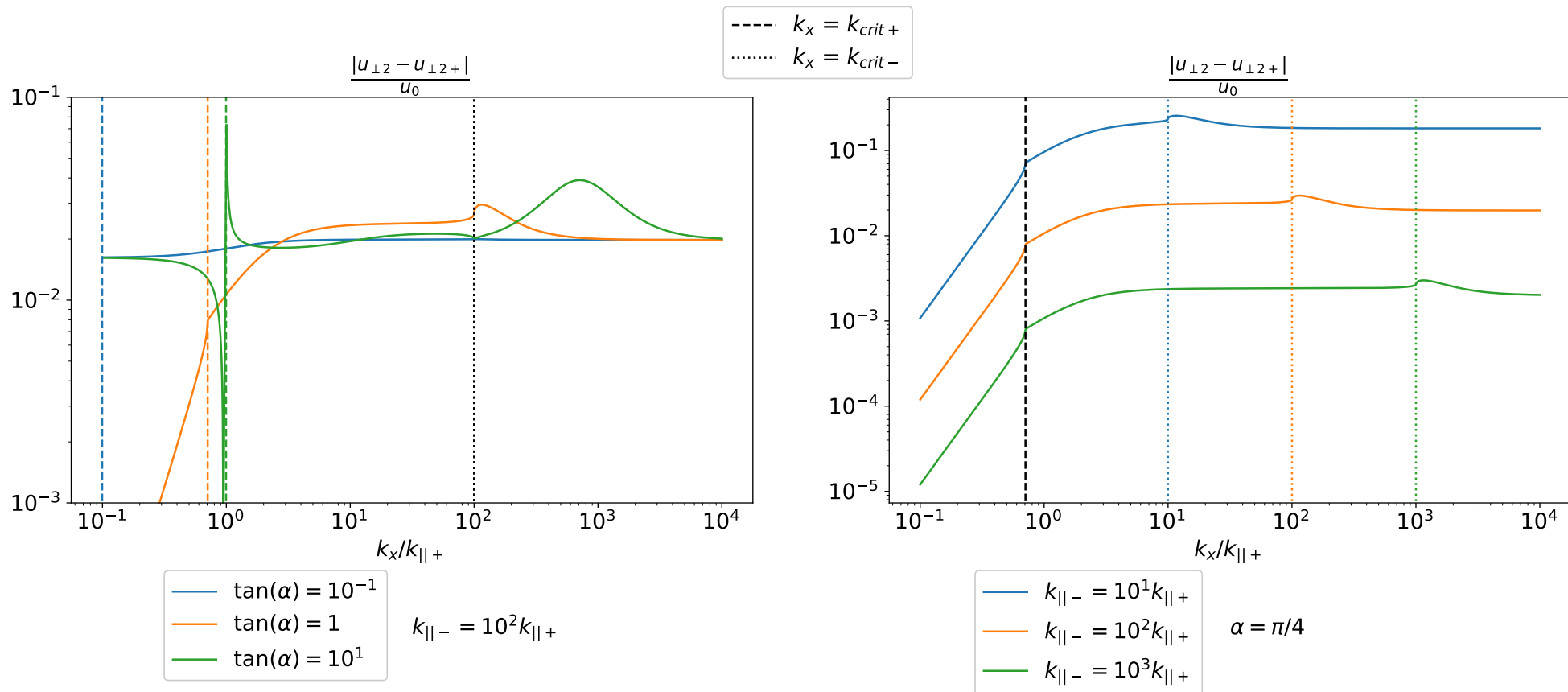


Chromosphere | Corona

Reflected fast wave error



Reflected Alfvén wave error



Summary + conclusions

- Alfvén waves couple to fast waves at the TR
- They change polarisation upon reflection
- Line-tied BC's are usually a good approximation
- However, they generate unphysically large BL's if:

$$k_x \gg k_{\parallel-}$$

Future work

- Investigate different incident polarisations
- Use incident fast waves
- Model an exponential density profile instead of piecewise constant
- Let $\rho = \rho(x, z)$

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